

# PETROLEUM ENGINEERING

## Mathematical Model for Flow of Pseudoplastic Fluids in Porous Media

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(Received 14/2/1993; accepted for publication 3/7/1993)

**Abstract.** A non-linear partial differential equation for the flow of all types of pseudoplastic fluids in porous media was derived. The rheological characteristics of the pseudoplastic fluids were described by the generalized approach developed by Metzner and Reed. A linear form of the equation developed was derived and solved. Analytical solutions were obtained for a constant injection rate into an infinite reservoir.

The validity of the present analysis was verified by comparing the pressures calculated from the developed equation with published data. The equation deduced provides an accurate prediction for pressure distributions.

### Nomenclature

- C = fluid compressibility, Pa<sup>-1</sup>
- C<sub>t</sub> = system compressibility, Pa<sup>-1</sup>
- D<sub>p</sub> = particle diameter, m
- h = formation thickness, m
- K = Consistency coefficient, Pa.S<sup>n</sup>
- k<sub>r</sub> = permeability in the radial (horizontal) direction, m<sup>2</sup>
- K' = Metzner and Reed parameter, Pa.s<sup>n'</sup>
- n' = Metzner and Reed parameter
- P = pressure, Pa
- P<sub>D</sub> = non-Newtonian dimensionless pressure
- P<sub>i</sub> = initial pressure, Pa

$P_{wf}$  = flowing bottomhole pressure, Pa

$\Delta P$  = Pressure difference,  $p_{wf} - p_i$ , Pa

$q$  = flow rate,  $m^3/s$

$r$  = radial distance, m

$r_D$  = dimensionless radial distance

$r_w$  = wellbore radius, m

$t$  = time, s

$t_D$  = non-Newtonian dimensionless time

$u$  = superficial velocity (volumetric rate of flow per unit cross-sectional area), m/s

$u_r$  = superficial velocity in the radial (horizontal) direction, m/s

$V$  = pore velocity, m/s

$\dot{\gamma}$  = shear rate,  $s^{-1}$

$\Gamma(x)$  = Gamma function.

$\mu_{eff}$  = effective viscosity (a viscosity level parameter for pseudoplastic fluids),  $Pa \cdot s^n \cdot m^{1-n}$

$\lambda$  = group defined by Eq. 12.

$\phi$  = porosity, fraction

$\tau$  = shear stress, Pa

### Subscripts

D = dimensionless

r = radial

w = wellbore

## Introduction

Non-Newtonian fluids such as polymer solutions, microemulsions and macroemulsions are injected into reservoirs to control mobility and enhance oil recovery. These fluids necessitate a basic understanding of their flow behavior through porous media [1;2]. Several studies of the rheology of non-Newtonian fluids in porous media have been reported in literature [3-17]. Most of these studies have mathematically described the flow of pseudoplastic fluids by either power-law or Ellis model [3-14]. Pascal [15;16], and Al-Fariss and Pinder [17] investigated the flow of yield-fluids through porous media. Thus, most of the equations developed by previous investigators for pseudoplastic fluids were limited to two models: power-law and Ellis models. These models are two of the twenty known rheological models of pseudoplastic fluids [18]. Thus, the equations developed by previous investigators cannot be

used to predict the pressures distribution in reservoirs for all types of pseudoplastic fluids.

Therefore, the present analysis was carried out to develop a general equation for all pseudoplastic fluids. The rheological characteristics of those fluids can be described by the generalized flow shear equation developed by Metzner and Reed [19].

$$\tau_w = K' (8V/D)^{n'} \quad (1)$$

### Rheological Model for Pseudoplastic Fluids

The porous media can be regarded as a conduit with a complicated cross section with mean hydraulic radius  $R_h$  equal to  $R/2$  [10]. The superficial velocity ( $u$ ) can be related to the average pore velocity ( $V$ ) by the following relation

$$u = \phi V \quad (2)$$

The wall shear stress in a porous medium is given by Bird, et al. [20]

$$\tau_w = \frac{2}{25} \frac{D_p \phi}{(1 - \phi)} \frac{\Delta P}{L} \quad (3)$$

The permeability of the porous medium can be defined from Blake-Kozeny Equation [20].

$$K_r = \frac{D_p^2 \phi^3}{150 (1 - \phi)^2} \quad (4)$$

The term 150 is an average value accepted by many investigators since its first use by Ergun [12]. Substituting Eqs. 2,3 and 4 in Eq. 1, the superficial velocity is expressed as

$$u = \left( \frac{K_r}{\mu_{\text{eff}}} \frac{\Delta P}{L} \right)^{1/n'} \quad (5)$$

where  $\mu_{\text{eff}}$  is given by

$$\mu_{\text{eff}} = 12^{n'-1} K' (150 \phi K_r)^{(1-n')/2} \quad (6)$$

For power-law fluids, the values of  $n'$  and  $K'$  are equal to  $n$  and  $K [(1+3n)/4n]^n$  respectively. Hence Eq. 6 becomes

$$\mu_{\text{eff}} = \frac{K}{12} \left( \frac{3+9n}{n} \right)^n (150 K_r \phi)^{(1-n)/2} \quad (7)$$

Eq. 7 is the same equation obtained by previous investigators [2,11,13] for power-law fluids. In radial coordinates, Eq. 5, can be written as

$$u_r^{n'} = \left( \frac{K_r}{\mu_{\text{eff}}} \frac{\partial P}{\partial r} \right) \quad (8)$$

Eq. 8 can be combined with the continuity equation and an equation of state to derive a partial differential equation for flow of all types of pseudoplastic fluids in porous media.

### Theoretical Analysis

A partial differential equation that models the flow of pseudoplastic fluids through a porous medium will be derive under the following conditions.

- 1- The porous medium has a uniform thickness and is homogeneous.
- 2- Permeability is constant.
- 3- The compressibility of the fluid and pressure gradient are small.
- 4- Gravity effect is negligible.
- 5- The flow of the fluid is isothermal.

The continuity equation for radial flow in a porous medium is given by Ikoku [13] and Desouky [18].

$$\frac{1}{r} \frac{\partial}{\partial r} (r \rho u_r) = - \frac{\partial}{\partial t} (\phi \rho) \quad (9)$$

The equation of state may be written as

$$\rho = \rho_0 e^{C(P-P_0)} \quad (10)$$

Substituting Eqs. 8 and 10 in Eq. 9, and applying the conditions of items (1) through (5), the following equation is obtained.

$$\frac{\partial^2 P}{\partial r^2} + \frac{n'}{r} \frac{\partial P}{\partial r} = \lambda \frac{\partial P}{\partial t} r^{1-n'} \quad (11)$$

where

$$\lambda = 12^{n'-1} n' \phi C_t K' (q/2\pi h)^{n'-1} (150 \phi K_r)^{\frac{1-n'}{2}} / K_r \quad (12)$$

$$C_t = C + \frac{1}{\phi} \frac{\partial \phi}{\partial \rho} \quad (13)$$

For power-law fluids, where  $n' = n$  and  $K' [(1+3n)/4n]^n$ , Eq. 11 reduces to the flow equation developed by Ikkoku [2].

#### Dimensional Analysis

It might be advisable to express Eq. 11 in dimensionless form. The following definitions will be used.

$$1) \quad P_D = (P - P_i) / \lambda_1 \quad (\text{Dimensionless pressure})$$

where

$$\lambda_1 = \lambda q r_w^{1-n'} / (2\pi h C_t \phi n')$$

$$2) \quad r_D = r / r_w \quad (\text{Dimensionless radius})$$

$$3) \quad t_D = \frac{t}{\lambda r_w^{3-n'}} \quad (\text{Dimensionless time})$$

$$4) \quad \frac{\partial P}{\partial r} = \frac{\partial P_D}{\partial r_D} \cdot \frac{\partial r_D}{\partial r} \lambda_1$$

$$5) \quad \frac{\partial^2 P}{\partial r^2} = \frac{\partial^2 P_D}{\partial r_D^2} \cdot \frac{\partial r_D^2}{\partial r^2} \lambda_1$$

$$6- \quad \frac{\partial r_D}{\partial r} = \frac{1}{r_w}$$

$$7- \quad \frac{\partial P}{\partial t} = \frac{\partial P_D}{\partial t_D} \cdot \frac{\partial t_D}{\partial t} \lambda_1$$

$$8- \quad \frac{\partial t_D}{\partial t} = \frac{1}{\lambda r_w^{3-n'}}$$

Substituting the equations defined by items 1 through 8 in Eq. 11, the following linear partial differential equation is obtained.

$$\frac{\partial^2 P_D}{\partial r_D^2} + \frac{n'}{r_D} \frac{\partial P_D}{\partial r_D} = r_D^{1-n'} \frac{\partial P_D}{\partial t_D} \quad (14)$$

#### Analytical Solution

Equation 14 can be solved for the case of constant injection rate of pseudoplastic fluids into an infinitely large reservoir. The initial and boundary conditions are:

$$(a) \quad P_D(r_D, 0) = 0.0$$

$$(b) \quad (\partial P_D / \partial r_D)_{r_D} = -1.0, \quad \text{for } t_D > 0$$

$$(c) \quad P_D(r_D, t_D) \rightarrow 0, \text{ as } r_D \rightarrow \infty$$

$$(d) \quad \text{at the wellbore, } r_D = 1$$

Using the Laplace transform method [22], applying the conditions of items a through d, and deriving the inversion of the resulting equation for large times, the following approximate analytical inversion is obtained.

$$P_D = \frac{3-n' \frac{2(1-n')}{(3-n')} t_D^{\frac{1-n'}{3-n'}}}{(1-n') \Gamma\left(\frac{2}{3-n'}\right)} \quad (15)$$

In dimensional variables, Eq. 15 becomes

$$P_{wt} = P_i + \left( \frac{q}{2\pi h} \right)^{n'} \frac{12^{n'-1} K'}{K_r} \frac{(150 \phi K_r)^{(1-n')/2}}{(1-n')} \left[ (3-n')^2 t/\lambda \right]^{(1-n')(3-n')} / \Gamma \left( \frac{2}{3-n'} \right) \quad (16)$$

Equation 16 represents the flow of any pseudoplastic fluid in a porous medium under the conditions of constant injection rate and an infinitely large reservoir. A graph of  $\log \Delta P$  versus  $\log t$  yields a straight line of slope  $(1-n')/(3-n')$ , which can be solved for the value of  $n'$ . The intercept at  $t = 1$ , may be solved for the value of  $K'$  if  $\phi$ ,  $K_r$  and  $C_i$  are known. In case of knowing the rheological properties of the injected fluid, the transient flow behavior of that fluid can be determined.

#### Determining the Values of $n'$ and $K'$

Values of  $n'$  and  $K'$  can be estimated from the measured data of shear stress  $\tau$  and shear rate  $\dot{\gamma}$  using the following Rabinowitsch and Mooney equation (15).

$$\frac{8V}{D} = \frac{4}{\tau_w^3} \int_0^{\tau_w} \tau^2 \dot{\gamma} d\tau \quad (17a)$$

where  $n'$  and  $K'$  are defined as:

$$n' = \frac{\ln [ d(\tau_w) ]}{\ln [ d(8V/D) ]} \quad (18)$$

$$K' = \tau_w / (8V/D)^{n'} \quad (19)$$

In the present analysis, the procedure followed to integrate Eq. 17a is the modified Gauss four-points method suggested by Desouky and El-Emam [23] where Eq. 17a is written as:

$$\frac{8V}{D} = 4 I / \tau_w^2 \quad (17b)$$

where  $I$  is given by

$$I = 0.173 [(\tau^2 \dot{\gamma})_1 + (\tau^2 \dot{\gamma})_4] + 0.326 [(\tau^2 \dot{\gamma})_2 + (\tau^2 \dot{\gamma})_3]$$

The procedure for determining  $8V/D$  as a function of shear stress and shear rate is outlined as follows:

- 1) Square the values of shear stress ( $\tau$ ) and multiply them by the shear rate ( $\dot{\gamma}$ ).
- 2) Make a suitable correlation relating  $\log(\tau^2 \dot{\gamma})$  and  $\log(\tau)$ .
- 3) At the first experimental data point of  $\tau_1$ , determine the values of the following four points of  $(\tau^2 \dot{\gamma})$  from the equations.

$$(\tau^2 \dot{\gamma})_1 \text{ is estimated at } \tau = 0.06945\tau_1$$

$$(\tau^2 \dot{\gamma})_2 \text{ is estimated at } \tau = 0.33\tau_1$$

$$(\tau^2 \dot{\gamma})_3 \text{ is estimated at } \tau = 0.67\tau_1$$

$$(\tau^2 \dot{\gamma})_4 \text{ is estimated at } \tau = 0.93055\tau_1$$

- 4) Assume the calculated  $\tau_1$  equal to  $(\tau_w)_1$ , then calculate the first value of  $(8V/D)$  from Eq. 17b. The calculated values of  $(8V/D)$ , and  $(\tau_w)_1$  are corresponding to  $\tau_1$  and  $\dot{\gamma}_1$ .
- 5) Repeat steps (3) and (4) for the remaining experimented data points of  $\tau$  and  $\dot{\gamma}$  to obtain the corresponding converted values of  $8V/D$  and  $\tau_w$ .

### Comparison with the Published Data

To verify the validity of the present analysis, a comparison was made between the pressures calculated from Eq. 16 and the simulated data (2, 24). A biopolymer injection test was performed at a test site during various periods of the injection history. The measured pressures are plotted against time in Fig. 1. The average porosity, thickness and flow rate were found to be 0.15, 4.877 m, and  $1.84 \times 10^{-4} \text{ m}^3/\text{s}$  respectively. The rheological properties of the injected polymer are plotted in Fig. 2. Equations 17 through 19 were used to calculate the values of  $n'$  and  $K'$ . These values are 0.323 for  $n'$  and 0.021 Pa.s $^{n'}$  for  $K'$ . Equation 16 shows that the slope of a graph ( $\log \Delta P - \log t$ ) is given by

$$\text{Slope} = \frac{1 - n'}{3 - n'} \quad (20)$$

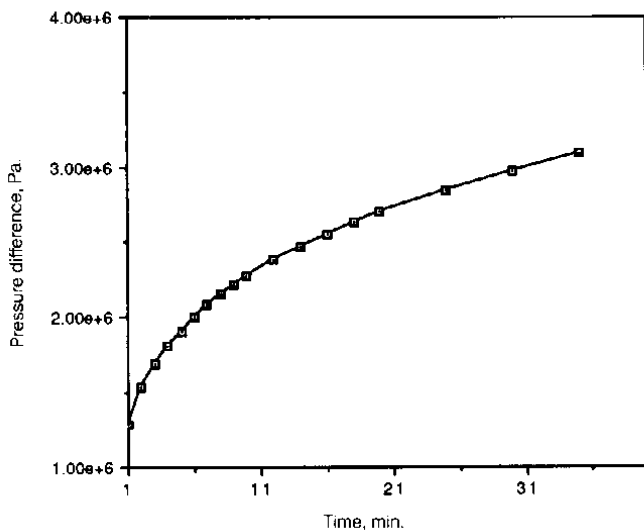


Fig. 1. Pressure difference vs. time for a polymer injection test.

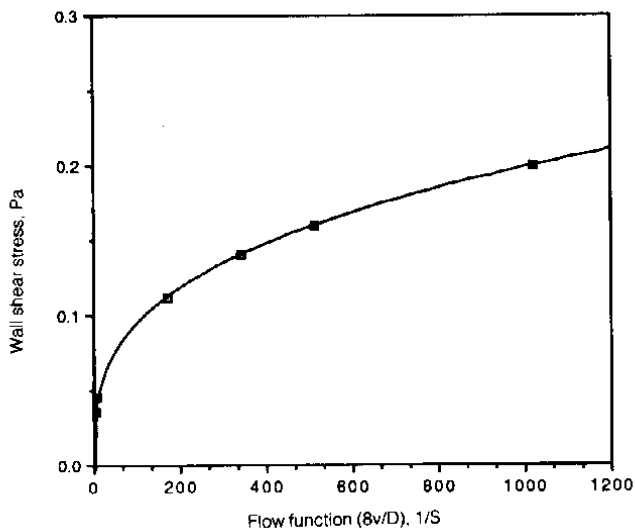


Fig. 2. Rheological data and resulting flow curve.

The value of the slope can be obtained by plotting the measured pressures versus time on log-log scale as shown in Fig. 3. This figure shows that the value of the slope is equal to 0.25. The value of  $n'$  was calculated from Eq. 20 and was found to be 0.333. This indicates that the value of  $n'$  determined from Fig. 3 is consistent with that obtained from rheological data given in Fig. 2. The pressures calculated from Eq. 16 are plotted against the measured ones in Fig. 4. This figure shows that an excellent agreement exists between the pressures calculated from Eq. 16 and the measured ones with an average relative error of 0.32%.

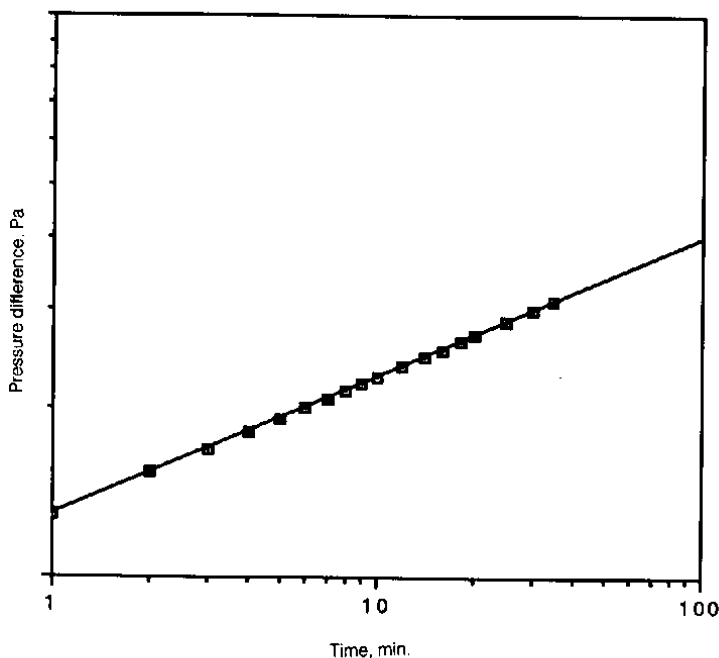


Fig. 3. Logarithm of pressure difference vs. logarithm of time for a polymer injection test.

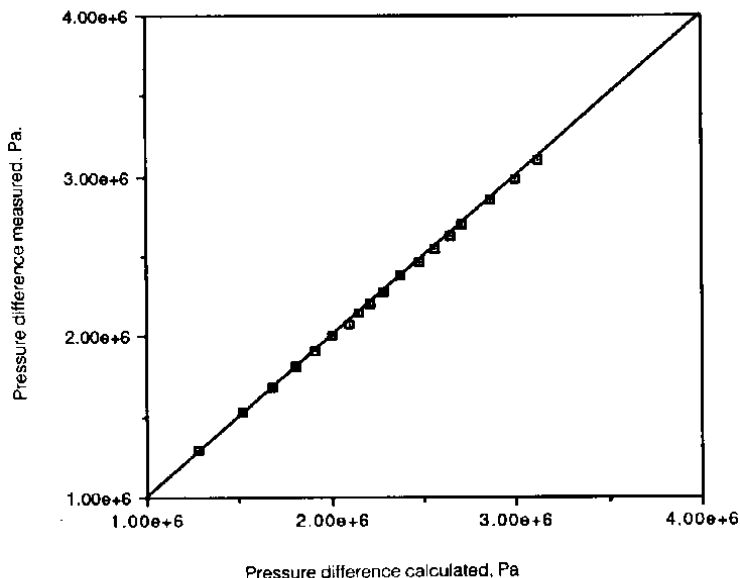


Fig. 4. Cross plot for the measured and calculated pressure drop.

### Conclusion

Flow of pseudoplastic fluids in porous media is expressed by a simple mathematical formula, Eq. 16. It can be used for any type of pseudoplastic fluids irrespective of the rheological model which described the relationship between shear stress and shear rate. The equation developed ables to predict the pressure drop in a porous medium due to flow of a pseudoplastic fluid with an average relative error of 0.32%.

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## نموذج رياضي لسريان السوائل اللدائنية الكاذبة في الأوساط المسامية

عادل محمد حميدة

قسم هندسة النفط، كلية الهندسة، جامعة الملك سعود، ص. ب. ٨٠٠،

الرياض ١١٤٣١، المملكة العربية السعودية

(استلم في ١٤/٢/١٩٩٣ م؛ قُبل للنشر في ٣/٧/١٩٩٣ م)

ملخص البحث. في هذا البحث استنبطت معادلة تفاضلية جزئية غير خطية لسريان السوائل اللانيوتونية (اللدائنية الكاذبة) في الأوساط المسامية. عبّر عن الخواص الريولوجية لهذه السوائل بالطريقة العامة لمزج وريد. واقترحت حلول تحليلية خطية للمعادلة المستنبطة عند معدّل حقن ثابت داخل المكنم اللاحدودي.

تم اختبار دقة المعادلة المقترحة وذلك بمقارنة النتائج المحسوبة منها بالبيانات المنشورة.